Chemical ordering in PtNi nanocrystals

Okkyun Seo^a, Ji Yeon Lee^a, Jae Myung Kim^a, Jin-Woo Kim^a, Hyon Chol Kang^b, Jinwook Chung^c, Do Young Noh^{a,*}

^aDepartment of Physics and Photon Science & School of Materials Science and Engineering, Gwangju Institute of Science and Technology, Gwangju 61005, Korea

^bDepartment of Advanced Materials Engineering, Chosun Univ., Gwangju 61542, Korea

^cDepartment of Physics, Pohang Univ. of Science and Technology, Pohang 37673, Korea

Abstract

We investigated the chemical ordering in PtNi nanocrystals fabricated on sapphire substrate using *in-situ* synchrotron x-ray scattering. Nanocrystals with composition close to 1:1 were ordered in the tetragonal L1₀ structure at low temperatures. The transition to disordered FCC structure occurred at around 640 °C and substantial hysteresis of about 50 K was observed. Nanocrystals of smaller sizes fabricated under the same conditions were Ni rich and ordered into Cu₃Au type L1₂ structure. Significantly higher degree of chemical ordering was observed in L1₂ structure than in L1₀ structure.

Keywords: PtNi alloy nanocrystals, chemical ordering, order-disorder transition, L1₀ & L1₂ structure

1. Introduction

Recently, Pt based bimetallic alloy nanoparticles attracted much attention for applications in fuel cell electrodes and high density magnetic storage devices[1, 2, 3, 4] since the catalytic activity can be optimized and the amount of Pt can be reduced by alloying with less expensive metals. Among the Pt bimetallic alloys for fuel cell applications, PtNi nanocrystals have been extensively investigated due to their excellent oxygen reduction reaction efficiency (ORR) surpassing Pt. For example, the ORR of Pt₃Ni nanoparticles was reported to be much higher than Pt[5, 6]. Pt-Ni core-shell structures have been synthesized and shown to enhanced the catalytic activity for ORR[7]. PtNi nanoparticles exhibit only weak ferromagnetic or paramagnetic behavior different from other Pt bimetallic alloys such as PtCo and PtFe, and they are mostly considered for the catalytic applications rather than magnetic applications. Tunability in size, shape, alloy composition, and elemental distribution of Pt alloy nanocrystals provides a great opportunity to optimize desired specific properties for specific applications.

*Corresponding author

Email address: dynoh@gist.ac.kr (Do Young Noh)

Bulk Pt based bimetallic alloys are known to exhibit chemical ordering below order-disorder transition temperature (T_C) in tetragonal L1₀ or cubic L1₂ structure depending on the alloy composition [8, 9]. Tetragonal L1₀ structure is preferred when the compositions of Pt and alloying metals are similar. Above T_C , the alloys are chemically disordered and alloying elements are randomly distributed on FCC lattice sites. Structural behaviors including melting and ordering of Pt bimetallic alloys in a nanocrystal form might be different from bulk due to the effects of finite size. It has been reported that Pt alloy nanocrystals in a chemically ordered structure exhibit distinct magnetic properties and superior catalytic activities[10, 11, 12, 13]. Theoretical studies predicted that the nature of chemical ordering changes as the particle size decreases to a few nm scale [14, 15]. Understanding the role of surface atoms and the interplay between the morphology and atomic ordering is important in nanocrystals. [16]. Recent Monte Carlo and molecular dynamics study showed that the L1₀ structure in PtNi nanoparticles is very stable and short range order persists to very high temperatures[17].

Experimental studies on structure and chemical ordering in alloy nanocrystals such as order-disorder transition and phase segregation have been difficult due to their small sizes and complicated behaviors. Nanocrystals in powder form tends to agglomerate at high temperatures and easily oxidized under non vacuum environment. Distinguishing equilibrium thermodynamic behavior from kinetically limited behavior becomes difficult as the size of particles approach nanometer scale. Due to these reasons, experimental data on detailed structural properties including atomic ordering in alloy nano crystals are still lacking.

In this report, we investigated the atomic ordering of PtNi nanocrystals aligned on a sapphire substrate using the *in-situ* synchrotron x-ray scattering technique. Nanocrystals with size above 150 nm exhibited a transition from the ordered L1₀ structure to disordered FCC structure around 650 °C. We observed a temperature hysteresis of about 50 K in the order parameter and lattice structure across the transition, which was not reported previously. Nanocrystals of smaller sizes were Ni rich and exhibited much higher degree of ordering. The crystal structure of the ordered phase was cubic L1₂ rather than the tetragonal L1₀.

2. Experimental Details

The PtNi alloy nanocrystals were prepared by applying the solid-state dewetting process on Pt-Ni bilayer films which were grown on a sapphire (0001) substrate by electron beam evaporation at room temperature. The thickness ratio of each layer in the Pt-Ni bilayer were maintained to 1:1, while each layer thickness was varied to obtain PtNi alloy nano crystals of various sizes. In this experiment, we report the results from Pt-Ni bilayer films with layer thickness of 1, 2.5, and 5 nm. PtNi alloy nanocrystals were obtained by rapid thermal annealing (RTA) of the Pt-Ni bilayer films to 1050°C for 10 minutes. The ramping rate of RTA was about 200 °C/min. The bilayer films were transformed into

Figure 1: (a) SEM image of PtNi nanocrystals obtained by annealing a $\text{Pt}_{5nm}\text{-Ni}_{5nm}$ bilayer film. (b) X-ray diffraction profile of the PtNi (111) Bragg peak. The sharp rise in the left side is the tail of the substrate sapphire (0006) Bragg peak. The line is the result of a fit to estimate the peak width.

completely mixed disordered alloy nanocrystals through solid-state dewetting process during annealing[18].

To transform the nanocrystals into chemically ordered structure, we annealed the samples at 530 °C below the expected order-disorder transition temperature for 50 hours. The Pt ratio in the nanocrystals, estimated from the unit cell size by applying the Vegard's law, varied from about 0.54 to 0.37 depending on the crystal size. The crystal orientation of the PtNi nanocrystals thus obtained were well aligned to the substrate sapphire orientation, which facilitated the structural analysis using x-ray diffraction significantly. The $\langle 111 \rangle$ axis of the PtNi nanocrystals were well aligned to the sapphire $\langle 0001 \rangle$ direction, while the in-plane $\langle 11\bar{2} \rangle$ were parallel to the sapphire $\langle 11\bar{2}0 \rangle$ direction.

The lateral particle size and size distribution were investigated by scanning electron microscopy (SEM) as illustrated in Fig. 1(a) which shows the PtNi nanocrystals obtained by annealing a $\text{Pt}_{5nm}\text{-Ni}_{5nm}$ film. The diameter of the nanocrystals was 160 ± 50 nm. The particle size in the vertical direction, estimated from the width of the (111) Bragg peak shown in Fig. 1(b) using Scherrer's equation, was about 80 nm, which is about half the lateral size. As the size of nanocrystals decreases below 200 nm, the crystal domain size estimated by the width of the Bragg peaks approaches the physical height of nanocrystals. Typically the vertical height of the nanocrystals supported on sapphire substrates is about half the lateral size[18].

We carried out the *in-situ* synchrotron x-ray scattering experiments at 5D GIST beam-line of Pohang Light Source (PLS) in Korea and at beamline X20 at National Synchrotron Light Source(NSLS) in USA. The x-ray energy was fixed at 10 keV by using a double-bounce Si (111) monochromator and focused by an x-ray mirror. The samples were maintained in a temperature controlled chamber with an x-ray window maintained at the 3×10^{-5} Torr of vacuum. The chemical ordering was investigated using the superlattice (100) peak which is forbidden in the disordered FCC structure. To investigate the crystal structure,

Figure 2: *In-situ* x-ray diffraction profile of the superlattice (001) reflection (a) and the fundamental (200) Bragg reflection (b) obtained from the PtNi nanocrystals at various temperatures across the order-disorder transition.

we measured the fundamental (200) peak as well as the (111) Bragg peak in the substrate normal direction.

90 The order-disorder transition was investigated by measuring the (100) and the (200) peak as we vary the sample temperature across the transition temperature T_C . X-ray diffraction profiles were measured at each temperature after waiting for 30 min to establish equilibrium. Starting from the ordered structure at room temperature, the measurement was done first as the sample was
 95 heated to across T_C (heating run). The measurement was repeated during cooling the sample below T_C (cooling run). The reproducibility was confirmed by performing a second heating run following the cooling run.

3. Results and Discussion

3.1. Order-Disorder transition in PtNi nanocrystals in $L1_0$ structure

100 We first present the order-disorder transition investigated on the PtNi nanocrystals obtained by annealing a $\text{Pt}_{5nm}\text{-Ni}_{5nm}$ film (Sample #1) whose average lat-

eral (vertical) size was about 160 (80) nm. The Pt composition estimated from the unit cell volume was about 0.54. As signified by the well-defined sharp superlattice (001) peak in the longitudinal direction measured at room temperature illustrated in Fig. 2(a), the nanocrystals showed high degree of chemical ordering. The size of the chemically ordered domain or the correlation length of the chemical ordering estimated from the half width at the half maximum (HWHM) of the peak, was about 21 nm. We also measured the fundamental (200) Bragg peak, shown in Fig. 2(b), which occurred at a momentum transfer value smaller than the twice the (001) peak position. This indicates that the unit cell was elongated along the [100] direction and the ordered structure was the tetragonal L1₀. Together with the position of the (111) Bragg peak shown in Fig. 1(b), we obtained the unit cell dimensions in all direction. The fundamental peak was sharper than the superlattice peak and the crystal domain size estimated from the HWHM of the fundamental peak was about 60 nm. We conclude that one crystal domain included two to four ordered domains separated by domain boundaries.

The order-disorder transition occurred at around 640 °C as illustrated in the data shown in Fig. 2. As the sample temperature was increased to 600 °C, the (001) peak becomes a little sharper and shifted towards the low-Q direction indicating that the lattice was expanded and chemical ordering was improved. This improvement was due to the annealing effect, although the change was minimal. At 640 °C, however, the (001) superlattice peak intensity suddenly decreased and the fundamental (200) peak shifted to a higher-Q value with increased intensity. This indicates that the chemical ordering was lost and the crystal structure was changed. Above T_C , the crystal structure was found to be cubic FCC as confirmed by the lattice constants evaluated using the peak positions of the (200) and the (111) Bragg peak.

The order-disorder transition in the PtNi nanocrystals was a first order transition with a large temperature hysteresis. Figure 3 illustrates the behavior of the order parameter evaluated from the superlattice (001) peak intensity and the a - and c -axis lattice parameter evaluated from the fundamental (200) and (111) peak positions.

The order parameter was defined as the the sum of the fraction of Pt sites occupied by Pt atoms and the fraction of Ni sites occupied by Ni atoms[19]. It is zero for a completely random system and unity for completely ordered system. The order parameter S presented in Fig. 3 was evaluated from the ratio of the integrated intensity of the superlattice (001) peak, I_s and the fundamental (200) peak, I_f assuming that the CuAu type of chemical ordering. Specifically the following formula,

$$S^2 = \frac{I_s}{I_f} \left(\frac{L_f D_f}{L_s D_s} \right) \left(\frac{F_{Pt} + F_{Ni}}{F_{Pt} - F_{Ni}} \right)^2,$$

where $L_f(s)$ and $D_f(s)$ are Lorentz factor and Debye-Waller factor for the fundamental (superlattice) peaks, and F_{Pt} and F_{Ni} are atomic form factor of Pt and Ni respectively.

Figure 3: (a) Order parameter representing the chemical order in the PtNi nanocrystals evaluated from the superlattice (100) peak. (b) a - and c -axis constants evaluated from the fundamental (200) and (111) Bragg reflection. Broad hysteresis was observed between the heating and cooling run. The lines are guide-to-the-eyes.

Abrupt changes were observed both in the order parameter and in the lattice structure during heating and cooling run indicating that the transition was first-order. We observed a large hysteresis of about 50 K, that is, the transition during the cooling run occurred about 50 K below the transition temperature observed during the heating run. Although we cannot completely exclude the effect of kinetic limitations, the hysteresis observed in this work was rather large.

3.2. $L1_2$ ordered structure observed in Ni rich PtNi nanocrystals

PtNi nanocrystals obtained from thin $\text{Pt}_{2.5\text{nm}}\text{-Ni}_{2.5\text{nm}}$ (Sample #2) and $\text{Pt}_{1\text{nm}}\text{-Ni}_{1\text{nm}}$ (Sample #3) bilayer films exhibited chemical ordering different from Sample #1. The average lateral size of the nanoparticles in Sample #2 and #3 estimated from SEM images was about 120 nm. Shown in Fig. 4 are the superlattice (001) peak obtained from Sample #2 and #3 at room temperatures together with that from Sample #1 discussed in the previous section. First, we note that the (001) peaks of Sample #2 and #3 shifted towards the low- Q side significantly indicating that their unit cell was larger than Sample #1 in this direction. The fundamental (200) and (111) peak positions were also different from those of Sample #1 (data not shown). By evaluating the unit cell dimensions, we found that the crystal structure was close to cubic within 1 % as summarized in table 1. From the unit cell dimensions, we found that chemical ordered nanocrystals in Sample #2 and #3 were Ni-rich with reduced Pt ratio of about 0.38. Presumably, Pt atoms are segregated out to achieve a stable

Figure 4: Superlattice (001) Bragg peak of Sample #1 (a), #2(b), and #3 (c) measured at room temperature. The ordered L1₀ (a) and L1₂ (b) structures are illustrated as insets. The lines are guide-to-the-eyes.

ordered configuration during long annealing to obtain the ordered structure. From the data shown in table 1, we conclude that Sample #2 and #3 were ordered in Cu₃Au type cubic L1₂ structure rather than the tetragonal L1₀.

The superlattice (001) peak of Sample #2 and #3 was also much sharper and more intense indicating that the degree of chemical ordering was higher and the ordered domains were larger. The ordered domain size estimated from the HWHM of the superlattice peak was 40 and 31 nm respectively, which is larger than the domain size of Sample #3, 21 nm, in L1₀ structure. We think that the order domain size in these samples was close to the particle dimension in the <001> direction, and most nanocrystals were consisted with single ordered domain. The high degree of ordering and large domains were possible in the cubic L1₂ structure because there is no distinction between the *a*-type and *c*-type domains any more.

Table 1: Lattice constants, crystal structure, and composition

sample	a=b (Å)	c (Å)	structure	composition
#1	3.808	3.595	L1 ₀ (tetragonal)	Pt _{0.54} Ni _{0.46}
#2	3.671	3.684	L1 ₂ (cubic)	Pt _{0.39} Ni _{0.61}
#3	3.672	3.663	L1 ₂ (cubic)	Pt _{0.37} Ni _{0.63}

4. Conclusions

In this work, we investigated the chemical ordering in PtNi nanocrystals aligned on a sapphire substrate using *in-situ* x-ray scattering. Relatively large nanocrystals with composition close to 1:1 were ordered in tetragonal L1₀ crystal structure. The order-disorder transition was first order. A large thermal hysteresis of about 50 K was observed both in the order parameter and in the lattice structure, which, we think, is unusually large. Nanocrystals of smaller size became Ni rich during the ordering process and exhibited cubic L1₂ ordered structure of Cu₃Au type. The chemical ordering was superb and the ordered domain size was close to the size of nanocrystals.

As the particle size decreases below 100 nm, they tends to phase-separate and Ni rich particles would appear. When the composition deviates from 1:1 toward 3:1, we expect that PtNi nanocrystals exhibit the *Cu₃Au* type of ordering in L1₂ structure. Although the transition temperature of particles close to 100 nm is similar to the bulk value, higher transition temperature predicted in literatures might appear in much smaller particles. Further studies in smaller nanocrystals would be useful for complete understanding of the size and composition dependence of the chemical ordering in Pt bimetallic nano alloy crystals.

Acknowledgments

We acknowledge the support and fruitful discussions provided by Dr. J.J. Sweet at NSLS and Dr. O. Sakata at Spring-8. This work was supported by the National Research Foundation of Korea (NRF) grant funded by the Korean government (MSIP) through SRC (NRF-2015R1A5A1009962), and We also acknowledge the GSG Project through a grant provided by GIST in 2015 and the support by Institute for Basic Science (IBS).

References

- [1] C. Wang, G. Wang, D. van der Vilet, K.-C. Chang, N. M. Markovic, V. R. Stamenkovic, Phys. Chem. Chem. Phys. 12 (2010) 6933.
- [2] T. Toda, H. Igarashi, H. Uchida, M. Watanabe, J. Electrochem. Soc. 146 (1999) 3750.
- [3] F. Tournus, A. Tamion, N. Blanc, A. Hannour, L. Bardotti, B. Prevel, P. Ohresser, E. Bonet, T. Epicier, V. Dupuis, Phys. Rev. B 77 (2008) 144411.

- 205 [4] Y. Yu, W. Yang, X. Sun, W. Zhu, X.-Z. Li, D. J. Sellmyer, S. Sun, Nano Lett. 14 (2014) 2778.
- [5] H. Yang, W. Vogel, C. Lamy, N. Alonso-Vante, J. Phys. Chem. B 108 (2004) 11024.
- [6] V. R. Stamenkovic, B. Fowler, B. S. Mun, G. Wang, P. N. Ross, C. A. Lucas, N. M. Markovic, Science 315 (2007) 493.
- 210 [7] G. Wang, H. Wu, D. Wexler, H. Liu, O. Savadogo, J. Alloys and Compd 503 (2010) L1–L4.
- [8] C. E. Dahmani, M. C. Cadeville, J. M. Sanchez, J. L. Moran-Lopez, Phys. Rev. Lett. 55 (1985) 1208.
- [9] Y. L. Bouar, A. Loiseau, A. Finel, Phys. Rev. B 68 (2003) 224203.
- 215 [10] Y. J. Zhang, Y. T. Yang, Y. Liu, Y. X. Wang, L. L. Yang, M. B. Wei, H. G. Fan, H. J. Zhai, X. Y. Liu, Y. Q. Liu, N. N. Yang, Y. H. Wu, J. H. Yang, J. Phys. D:Appl. Phys. 44 (2011) 295003.
- [11] M. Chen, D. E. Nikles, J. Appl. Phys. 91 (2002) 8477.
- [12] L. Xiong, A. Manthiram, J. Electrochem. Soc. 152 (2005) A697.
- 220 [13] S. Koh, M. F. Toney, P. Strasser, Electrochim. Acta 52 (2007) 2765.
- [14] D. Alloyeau, C. Ricolleau, C. Mottet, T. Oikawa, C. Langlois, Y. L. Bouar, N. Braidy, A. Loiseau, Nat. Mater. 8 (2009) 940.
- [15] W. Qi, Y. Li, S. Xiong, S.-T. Lee, Small 6 (2010) 1996.
- [16] D. Cheng, D. C. W. Wang, S. Huang, J. Phys. Chem. C 112 (2008) 4855.
- 225 [17] D. Cheng, S. Yuan, R. Ferrando, J. Phys: Condens. Matter 25 (2013) 355008.
- [18] O. Seo, S. A. Oh, J. Y. Lee, S. S. Ha, J. M. Kim, J. W. Choi, J.-W. Kim, H. C. Kang, D. Y. Noh, J. Alloys Compd. submitted for publication.
- 230 [19] B. E. Warren, X-ray Diffraction, Addison-Wesley Publishing Co, Inc., Reading, Mass., USA, 1990.