Pair Spectrometer Hodoscope for Hall D at Jefferson Lab[☆]

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Abstract

We present the design of the pair spectrometer hodoscope fabricated at Jefferson Lab and installed in the experimental Hall D. The hodoscope consists of thin scintillator tiles; the light from each tile is collected using wave-length shifting fibers and detected using a Hamamatsu silicon photomultiplier. Light collection was measured using relativistic electrons produced in the tagger area of the experimental Hall B.

Keywords: Scintillator detector, Pair spectrometer, Silicon photomultiplier

1. Introduction

The new detector GlueX has been constructed in the experimental Hall D at Thomas Jefferson National Accelerator Facility. The main physics goal of the GlueX experiment [1] is to search for hybrid mesons with exotic quantum numbers using a beam of linearly polarized photons incident on a liquid hydrogen target. The linearly polarized photon beam in Hall D will be produced via the coherent bremsstrahlung process by 12 GeV electrons in a thin diamond radiator. Coherent bremsstrahlung has been successfully used to produce linearly polarized photon beams at various experimental facilities [2, 3]. Coherent radiation from a diamond crystal lattice results in sharp monochromatic peaks in the photon energy spectrum. The peak energy depends on the relative orientation of the electron beam direction and the crystal lattice, and can be

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adjusted in the experiment by rotating the diamond. The GlueX photon energy region of interest corresponding to the main coherent peak is between 8.4 GeV and 9.1 GeV. The fraction of linearly polarized photons can be increased by passing a photon beam through a 3.4 mm diameter collimator situated about 75 m downstream of the radiator. The collimator preferentially filters out photons produced via coherent bremsstrahlung, which have a substantially larger mean emission angle. The collimated photons propagate toward the GlueX target. The energy of a beam photon can be inferred by analyzing the momentum of the electron after it has radiated the photon, given that the primary beam energy is known. This so-called “tagged” electron is deflected in a 6 m long dipole magnet operated at a field of 1.8 T and is subsequently registered in the tagging microscope or broad-band hodoscope scintillator counters.

One of the key components of the Hall D photon beam line is the magnetic pair spectrometer, which is installed after the photon collimator in front of the GlueX detector. The spectrometer will reconstruct the energy of a beam photon by detecting the $e^\pm$ pair produced by the photon in a thin converter. The main purpose of the spectrometer is to measure the spectrum of the collimated photon beam and determine the fraction of linearly polarized photons in the coherent peak energy region. It will also monitor the photon beam flux and can be used for the energy calibration of the tagging hodoscope and microscope detectors. The description of the Hall D pair spectrometer will be presented in Section 2. The design of the pair spectrometer high-granularity hodoscope and light collection measurements from hodoscope scintillator tiles will be described in Section 3 and Section 4, respectively.

2. Pair Spectrometer

Layout of the Hall D pair spectrometer is presented in Fig. 1. Electron-positron pairs are created by beam photons inside a thin converter with a typical thickness ranging between $10^{-4} - 10^{-2}$ radiation lengths. The choice of the converter thickness depends on the photon beam flux. The maximum flux for GlueX physics runs is expected to be $5 \cdot 10^7$ γ /sec in the coherent peak energy region. Three converters with different thicknesses are installed in a movable fork, that can insert one of them into the photon

Figure 1: Schematic top view of the Hall D pair spectrometer. PS and PSC denote the hodoscope and coarse scintillator counters, respectively.

41 beam. Produced leptons are deflected in a 18D36 dipole magnet with an effective field
 42 length of about 0.94 m. The magnet was brought from Brookhaven National Labora-
 43 tory and was modified at Jefferson Lab by reducing the pole gap from 6 inches to 3
 44 inches. The magnet is operated at a nominal field of 1.8 T. A 1.5 meter long vacuum
 45 chamber is installed after the magnet. Electrons and positrons are registered in two
 46 layers of scintillator detectors: a high-granularity hodoscope and a set of coarse coun-
 47 ters, denoted in Fig. 1 as PS and PSC, respectively. The detectors are organized into
 48 two arms positioned symmetrically with respect to the photon beam line. Each detector
 49 arm covers a momentum range of $e^\pm$ between 3.0 GeV/c to 6.2 GeV/c, corresponding
 50 to reconstructed photon energies between 6 GeV and 12.4 GeV. Relatively large accep-
 51 tance of the hodoscope allows one to reconstruct photons with energies in the coherent
 52 peak energy region and also in the range near the beam end-point energy of 12 GeV.
 53 This can be used for the energy calibration of the hodoscope detectors.

54 The high-resolution hodoscope is used for precise measurements of $e^\pm$ momenta.
 55 The hodoscope consists of a set of thin scintillator tiles, stacked together. Momentum
 56 of the reconstructed lepton is related to the x -coordinate of each tile, where x is an axis
 57 perpendicular to the beam line. Hodoscope tiles will be calibrated in units of lepton

energy, the energy scale will be used hereafter in the text. The hodoscope momentum resolution was studied using a detailed Geant detector simulation. The hodoscope energy resolution varies between 12 MeV and 20 MeV for 6 GeV and 12 GeV reconstructed photons, respectively. Light from each scintillator tile is detected using a Hamamatsu silicon photomultiplier (SiPM). The detailed description of the hodoscope will be presented in Section 3.

Sixteen coarse scintillator counters, eight in each detector arm, are positioned about 40 cm behind the hodoscope. The counters are 4.4 cm wide in the direction perpendicular to the lepton trajectory, 2 cm thick, and 6 cm in height. Hamamatsu R6427-01 PMTs are used to detect the scintillation light. The counters are used to produce a pair spectrometer trigger by requiring a coincidence of hits in the two detector arms. They also help to reduce background originating from interactions of $e^\pm$ inside the magnet pole edges to the level below 1% by constraining $e^\pm$ trajectories. Counter rates depend on the converter thickness and photon beam flux. The maximum rate will not exceed 10 kHz per PSC counter during GlueX operation.

Signals from both the hodoscope and coarse counters are digitized using a twelve-bit multi-channel flash ADC operated at a sampling rate of 250 MHz [5]. The PSC counters are also instrumented with TDCs, with a channel width of 60 ps [6]. The time resolution is expected to be better than 200 ps. This resolution will allow one to distinguish the electron beam bunch where the bremsstrahlung photon is emitted and therefore relate hits from the pair spectrometer and tagging detectors originating from the same event.

3. Pair Spectrometer Hodoscope

3.1. Design and Fabrication

The pair spectrometer hodoscope consists of 145 rectangular tiles made of EJ-212 scintillator [4], stacked together as shown in Fig. 2. The tile height is 3 cm and the length along the electron path in scintillator is 1 cm. Tiles with two different widths, in the direction perpendicular to the electron trajectory, of 2 mm and 1 mm are used. Forty 1 mm wide tiles are instrumented in the detector high-energy region around 6

87 GeV in order to provide better energy resolution in this energy range. The energy bin
88 size of the hodoscope tiles varies depending on the lepton energy and constitutes about
89 13 MeV and 24 MeV for 3 GeV and 6 GeV leptons, respectively. Tiles are optically
90 isolated using 10 μm aluminized Mylar foil. This reflective foil also covers the bottom
91 of the tile assembly. The deflection angle of leptons in the dipole magnet depends on
92 the lepton momentum; the angular spread is about 5 degrees between 3 GeV and 6.2
93 GeV electrons. To account for the angular dependence, hodoscope tiles are organized
94 into 18 groups. Each group is tilted by about 0.3 degrees. This was done by attaching a
95 shim, a $\sim 50 \mu\text{m}$ thick adhesive strip, to one side of the aluminum foil after each group.

96 Light from a tile is collected using two 20 cm long 1 mm x 1 mm square double-
97 clad BCF-92 wave-length shifting (WLS) fibers, each of which is glued to the side of
98 the tile using BC 600 Optical Cement. A tile assembly with two WLS optical fibers
99 is shown in Fig. 3. The peak of the emission spectrum for EJ-212 scintillator occurs
100 at the wavelength of 423 nm, which couples well with the absorption spectrum of the
101 BCF-92 fiber. Light is subsequently reemitted inside the fiber in the green range with
102 an emission peak of 492 nm.

103 Collected light is transmitted to the end of the WLS fiber. A pair of fibers from each
104 tile is inserted into a hole in an aluminum mounting plate, as shown on the upper plot
105 of Fig. 2. An electronics circuit board containing 145 photo detectors is attached to
106 the mounting plate. Each photo sensor is coupled to two WLS fibers from a single tile.
107 The light detection is performed using Hamamatsu surface mount S10931-050P silicon
108 photomultipliers with an effective photosensitive area of 3 mm x 3 mm and a pixel size
109 of 50 μm . These sensors have a photon detection efficiency (PDE) larger than 20% at
110 a wavelength of 500 nm and a typical gain of about $7 \cdot 10^5$ [7]. The electronics board
111 with SiPMs is presented in Fig. 4. The SiPMs are arranged in two arrays of 3 x 35 and
112 5 x 8 sensors, which are connected to 2 mm and 1 mm tiles, respectively. SiPMs are
113 optically isolated using a plastic spacer.

114 3.2. *Electronics*

115 The hodoscope front end electronics consists of an amplifier and a SiPM bias volt-
116 age control circuit developed at Jefferson Lab. Signals from the SiPMs are amplified

117 using the amplifier with a gain of about a factor of 20. The amplifier is based on
118 commercially available devices (operational amplifiers) with 3 GHz bandwidth. Pulse
119 shaping is employed to compensate for the characteristically high SiPM capacitance
120 and package inductance. The impulse response shows rise and fall times of 3 ns and
121 with trans-impedance gain of 1 mV/ μ A. The SiPM operating bias voltage is about 73
122 V. The nominal bias setting is as specified by Hamamatsu (1 V over voltage) and fed
123 to the SiPM through a resistive network employing a thermistor and a linearizing resistor.
124 The hodoscope control electronics supply individually adjusted voltages to groups
125 of 5 SiPM channels; inside the group the voltage is adjusted among channels using
126 resistors. The thermistor senses the average temperature of closely packed SiPMs in
127 thermal equilibrium via a heat spreader PCB layout, thus forming a well controlled
128 loop. The commercially available bias power supply has very low noise characteristics
129 and is well regulated to less than 1 mV long term. The supply allows the user to
130 monitor and adjust the levels as needed and if required. The optimal bias setting will
131 be determined based on experimental conditions. Amplified SiPM signals are digitized
132 by flash ADCs positioned in two readout VXS crates. An example of the flash ADC
133 signal pulse obtained from a scintillator tile is shown in Fig. 5. The ADC sampling
134 time is 4 ns.

135 After fabrication, the performance of SiPMs and electronics were tested using a
136 picoseconds laser light pulser¹. The typical rise time of a SiPM pulse produced by a
137 laser pulser is about 6 ns and the corresponding fall time is 45 ns. The pulser trigger
138 output was used to initiate the flash ADC readout. SiPM amplitudes digitized by ADC
139 are integrated in a time window of 60 ns. A typical ADC spectrum is shown in Fig. 6.
140 Peaks on this plot correspond to the number of pixels fired in the SiPM. The position
141 of a single-pixel peak depends on the combined gain of a SiPM and electronics. The
142 single-pixel peak distribution measured for all hodoscope SiPMs is presented in Fig. 7.
143 The solid curve corresponds to the fit of this distribution to a Gaussian function. The
144 relative gain variation is found to be smaller than 1.5%.

¹Hamamatsu PLP-10 light pulser was used to test the hodoscope electronics.

145 4. Light Collection Studies

146 4.1. Beam Test

147 During the design phase of the hodoscope, we studied light collection from scin-
148 tillator tiles using a detector prototype that consisted of fourteen 2 cm x 1 cm x 1 mm
149 and five 2 cm x 1 cm x 2 mm tiles. Tiles were made from the same scintillator type
150 EJ-212 but had smaller height of 2 cm. Light collection was performed using a sin-
151 gle BCF-92 fiber, which was glued to one side of the tile using DYMAX UV optical
152 adhesive. Light was detected by Hamamatsu SiPMs S10362-11-025C and S10362-11-
153 050C, which have the pixel size of 25 μm and 50 μm , respectively. These SiPMs have
154 a smaller sensitive area of 1 mm x 1 mm. Prototypes of the SiPM amplifier and the
155 bias control circuit were used. The trigger was provided by a small plastic scintillator
156 mounted behind the tiles. The trigger scintillator was connected to a Philips XP 1911
157 photomultiplier tube via two optical fibers. The setup used in the beam test is shown in
158 Fig. 8. SiPM pulses were digitized using a flash ADC in a VXS crate.

159 The hodoscope prototype was positioned in the tagger area of the experimental Hall
160 B at Jefferson Lab [3] and was operated concurrently with the CLAS detector. The
161 CLAS used a secondary beam of photons produced in the tagger area by the Jefferson
162 Lab electron beam in a thin radiator via the bremsstrahlung process. A beam electron
163 that radiates a photon is deflected in the dipole tagger magnet and sent to the floor of
164 the experimental hall. The prototype was placed inside a light-tight box and positioned
165 on the floor of the tagger area. The hodoscope was oriented in such a way that it
166 faced normal to the path of electrons, i.e., electrons travel 1 cm through the tile. A
167 schematic view of the tagger area and the location of the prototype is shown in Fig. 9.
168 Dashed curves on this plot denote trajectories of tagged electrons with the fractional
169 energy k/E_0 , where k is the tagged electron energy and E_0 is the electron beam energy.
170 Data was taken with an electron beam energy of 2.54 GeV, corresponding to a tagged
171 electron energy of about 1.4 GeV.

The amount of light collected from a tile was measured in units of pixels fired in
the SiPM. The average number of pixels was estimated to be

$$N_{\text{pixels}} = A_{\text{peak}}/A_{\text{single}}, \quad (1)$$

172 where A_{peak} is the ADC value corresponding to the maximum of an ADC spectrum
 173 measured in the beam test and A_{single} is the ADC value of a single pixel peak. The
 174 position of the single pixel peak was calibrated using a LED pulser. Exact counting
 175 of the number of photons that induce pixel-breakdowns is complicated due to the op-
 176 tical cross-talk and after-pulses in the SiPM. The cross-talk is produced by secondary
 177 photons emitted in the pixel breakdown that trigger an additional avalanche in a neigh-
 178 boring pixel. The after-pulsing is believed to be produced by charge carriers that are
 179 released from the substrate after some delay from the avalanche resulting in consec-
 180 utive SiPM pulses. The cross-talk and after-pulse probabilities for the SiPMs used in
 181 the hodoscope are expected to be relatively small, less than 10% [8]. For the prototype
 182 operational bias voltage (about 1 V over) we expect on average about 0.2 pixels fire
 183 per incident photon. SiPM operational voltages for the final detector will be adjusted
 184 according to the experimental conditions and detector performance.

185 The beam test results are listed in Table 1. The amount of light collected from the
 186 1 mm and 2 mm wide tiles is found to be almost the same and corresponds to about
 187 29 pixels fired in the SiPM with the pixel size of 50 μm . This indicates that the light
 188 collection mainly depends on the surface of the WLS fiber glued to the tile, which is
 189 similar for 1 mm and 2 mm tiles. The fewer number of fired pixels measured with a
 190 25 μm SiPM is consistent with about a factor of two smaller photon detection efficiency
 191 of these SiPMs compared with 50 μm SiPMs. The smaller pixel fill factor (the ratio of
 192 the active area of a pixel to the entire area) accounts for the smaller PDE for the 25 μm
 193 S10362-11-025C SiPMs. We estimate that errors on light collection measurements are
 194 dominated by alignment of the WLS fiber with respect to the SiPM. The sensitive area
 195 of SiPMs used in the prototype is the same as the fiber cross section.

196 Light collection from the final hodoscope detector will be performed using two
 197 WLS fibers and SiPMs with a factor of 9 larger effective area. Therefore, we expect
 198 about a factor of two more light collected from the tile and a SiPM response corre-
 199 sponding to about 60 fired pixels. This amount of light should be sufficient to provide a
 200 high detection efficiency for $e^\pm$ pairs in the hodoscope. The hodoscope can be operated
 201 at relatively small ADC thresholds because a coincidence of hits in two detector arms

Tile width (mm)	1		2
SiPM pixel size (μm)	25	50	50
Light detected (pixels)	16.1 ± 2.1	28.7 ± 3.4	29.2 ± 3.6

Table 1: The amount of light in units of fired pixels seen by SiPM for various tiles used in the beam test.

is required.

4.2. Cosmic Setup

During the fabrication of scintillator tiles for the final hodoscope, we checked light collection using a cosmic ray setup presented in Fig. 10. Muon candidates were selected using two small trigger scintillator counters with an effective area of 6 mm x 3 cm and thickness of 2 mm. Counters were positioned perpendicular to each other and formed an overlap region of 6 mm x 6 mm. The distance between scintillators was about 20 cm. A hodoscope tile was placed between these counters, in such a way that cosmic particles go perpendicular to the tile, in the middle. The average muon path length in the scintillator corresponds to the tile width. A lead brick about 10 cm thick was position above the bottom scintillator to filter low-energy particles. Light detection was performed using a Hamamatsu S10931-050P SiPM instrumented with a single-channel amplifier. The amplifier gain was about a factor of ten smaller than that used in the hodoscope. A typical ADC spectrum obtained from the SiPM is shown on the top plot of Fig. 11. The average amount of light collected from 2 mm wide tiles corresponds to about 18 pixels fired in the SiPM. A Sr90 source was used to check quality of some selected fabricated tile assemblies. The SiPM spectrum obtained from a 2 mm wide tile using the radioactive source is shown on the bottom plot of Fig. 11 for comparison.

5. Summary

We have described the design and fabrication details of the pair spectrometer hodoscope, an array of thin scintillator tiles. Light from each tile is detected using a 3

mm x 3 mm Hamamatsu SiPM. A detector prototype was built to perform light collection studies using relativistic electrons produced in the experimental Hall B at Jefferson Lab. The amount of light for 1 mm and 2 mm tiles measured with the prototype corresponds to a SiPM response of about 29 fired pixels. At least a factor of two more light is expected for the final hodoscope for which the light collection is performed using two WLS fibers and tiles with larger size, i.e., larger light collection surface, are used. This amount of light is sufficient to provide a high detection efficiency of leptons by the hodoscope. Two arms of the hodoscope detector were commissioned and installed in the experimental Hall D.

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Figure 2: Schematic view of the pair spectrometer hodoscope (top). Light from scintillator tiles is collected using WLS fibers, which are inserted into holes in the mounting plate. Fabricated detector is shown in the bottom picture.

Figure 3: Scintillator tiles with two WLS fibers glued to sides of each tile: 3 cm x 1 cm x 1 mm tile (left) and 3 cm x 1 cm x 2 mm tile (right).

Figure 4: Electronics board with 145 silicon photomultipliers.

Figure 5: A typical flash ADC signal pulse obtained from a scintillator tile.

Figure 6: An example of the ADC spectrum obtained from SiPMs during electronics tests using a laser pulser. The first peak on the spectrum around ADC count zero corresponds to the ADC pedestal.

Figure 7: Distribution of the single-pixel peak measured by flash ADC for 290 hodoscope channels.

Figure 8: Top view of the hodoscope prototype with the trigger scintillator used in the beam test in experimental Hall B.

Figure 9: Hall B tagger system. Location of the hodoscope prototype under the Hall B tagger. Dashed curves denote trajectories of tagged electrons.

Figure 10: A cosmic ray setup.

Figure 11: ADC spectra for a 2 mm wide tile measured with cosmic muons (a) and a ^{90}Sr radioactive source (b).